A Survey of Large Molecules of Biological Interest toward Selected High Mass Star Forming Regions

A. Remijan^{1,2,3}, Y.-S. Shiao¹, D.N. Friedel¹, D.S. Meier¹, L.E. Snyder¹

ABSTRACT

We have surveyed three high mass Galactic star forming regions for interstellar methanol (CH₃OH), formic acid (HCOOH), acetic acid (CH₃COOH), methyl formate (HCOOCH₃), methyl cyanide (CH₃CN), and ethyl cyanide (CH₃CH₂CN) with the BIMA Array. From our observations, we have detected two new sources of interstellar HCOOH toward the hot core regions G19.61-0.23 and W75N. We have also made the first detections of CH₃CH₂CN and HCOOCH₃ toward G19.61-0.23. The relative HCOOH/HCOOCH₃ abundance ratio toward G19.61-0.23 is 0.18 which is comparable to the abundance ratios found by Liu and colleagues toward Sgr B2(N-LMH), Orion and W51 (\sim 0.10). We have made the first detection of HCOOCH₃ toward W75N. The relative HCOOH/HCOOCH₃ abundance ratio toward W75N is 0.26 which is more than twice as large as the abundance ratios found by Liu and colleagues. Furthermore, the hot core regions around W75N show a chemical differentiation between the O and N cores similar to what is seen toward the Orion Hot Core and Compact Ridge and W3(OH) and W3(H₂O). It is also apparent from our observations that the high mass star forming region G45.47+0.05 does not contain any compact hot molecular core and as a consequence its chemistry may be similar to cold dark clouds. Finally, the formation of CH₃COOH appears to favor HMCs with well mixed N and O, despite the fact that CH₃COOH does not contain a N atom. If proved to be true, this is an important constraint on CH₃COOH formation and possibly other structurally similar biomolecules.

Subject headings: ISM: abundances and molecules - ISM: individual (G19.61-0.23) - ISM: individual (G45.47+0.05)- ISM: individual (W75N)

¹Department of Astronomy, University of Illinois, Urbana, IL 61801 email: aremijan@astro.uiuc.edu, shaw2@astro.uiuc.edu, friedel@astro.uiuc.edu, meierd@astro.uiuc.edu, sny-der@astro.uiuc.edu

²Current Address: NASA Goddard Space Flight Center, Earth and Space Data Computing Division, Code 930, Greenbelt, MD 20771

³National Research Council Resident Research Associate

1. Introduction

Recently, there has been great interest in searching a variety of star forming regions (SFRs) for biologically interesting molecules (biomolecules) including formic acid (HCOOH) and acetic acid (CH₃COOH) because they share common structural elements with glycine (NH₂CH₂COOH), the simplest amino acid (Remijan et al. 2003; Liu et al. 2002, 2001). By measuring the abundance and distribution of these structurally similar species, we are helping to constrain the number of sources that may contain interstellar glycine. Furthermore, large molecules such as CH₃COOH and HCOOCH₃ may provide important clues to the formation chemistry of large biomolecules. CH₃COOH, HCOOCH₃ and glycolaldehyde (CH₂OHCHO) are all isomers. However, HCOOCH₃ and CH₃COOH are detectable in the high mass SFRs W51e2, G34.3+0.24 and Sgr B2(N-LMH) and show a very compact (\sim 2-10 arcsec²) distribution, while CH₂OHCHO has only been detected toward Sgr B2(N-LMH) and shows a distribution extended over one arcmin (Hollis et al. 2001). The different distributions demonstrate that these isomers form in different ways. The question remains as to why CH₃COOH and HCOOCH₃ have such a different distribution compared to CH₂OHCHO despite their identical atoms and atomic composition. By further investigating the sources where large biomolecules may be found, a better understanding of the chemistry under which these species are formed can be achieved.

To date, the only interferometric detections of HCOOH and CH₃COOH have been toward high-mass (>10 $\rm M_{\odot}$) SFRs (Remijan et al. 2003, 2002; Liu et al. 2001). Two features that identify regions of high-mass star formation are ultracompact (UC) HII regions and H₂O masers. Dense regions of hot molecular gas are often associated with these regions and these "hot molecular cores" (HMCs) have typical molecular hydrogen densities of $n \sim 10^{5-7}$ cm⁻³ and temperatures of 100 - 200 K. It is toward these regions in and surrounding HMCs where the bulk of large biomolecules are found. Therefore, to extend our database of sources that contain large biomolecules, we carried out observations toward three high-mass SFRs with the Berkeley-Illinois-Maryland Association (BIMA) Array⁴. Because arrays typically have better angular resolution than single element radio telescopes and have larger fields of view, they are more suitable for determining the size and position of compact emission sources. For species with compact distributions this property will also lead to higher intensity lines and larger beam averaged column densities.

Remijan et al. (2003) carried out an extensive survey of CH₃COOH toward both high mass and low mass SFRs. As a result of their survey, they detected a new source of interstellar

⁴Operated by the University of California, Berkeley, the University of Illinois, and the University of Maryland with support from the National Science Foundation.

CH₃COOH toward the high-mass HMC G34.3+0.2. Furthermore, that survey suggested hot molecular cores that have a mass range between 200 and 2000 M_☉ and do not show differentiation between O and N chemistries are sources of interstellar CH₃COOH. Also, the three known CH₃COOH sources are within 7 kpc of the Galactic center. Therefore, it appears galactocentric distance could be a factor in the abundance of CH₃COOH. Finally, Remijan et al. (2004) found that methyl cyanide (CH₃CN) is a better tracer of the large molecules CH₃CH₂CN and CH₃COOH than is HCOOCH₃ toward the HMCs W51 e1 and e2. In order to investigate these hypotheses, we observed CH₃CN, CH₃CH₂CN, CH₃OH, HCOOH, CH₃COOH and HCOOCH₃ toward a sample of massive SFRs selected based on the above criteria.

2. Observations

The large molecule survey took place from 1999 October to 2003 November utilizing the B and C configurations of the BIMA Array. The minimum baseline for these observations was \sim 60 m, and the maximum baseline was \sim 150 m, providing angular resolutions of \sim 2"-13". Table 1 lists the parameters of the combined datasets of our survey⁵. To include all the data from a source with multiple tracks or in multiple arrays, the data were inverted in u-v space. The resulting synthesized beamsize was determined by the region of heaviest u-v coverage. Column (1) lists the sources observed. Column (2) lists the sources used to calibrate the antenna based gains. Saturn and Uranus were used as the flux density calibrators for these observations. The absolute amplitude calibration of these sources is accurate to within \sim 20%. Column (3) lists the corresponding beamsize; Column (4) lists the rest frequency of each observation and finally, column (5) lists the approximate 1 σ channel rms noise level in the CH₃COOH or CH₃CN windows of the combined datasets. The passbands were calibrated automatically during data acquisition⁶. In the past, this method has been quite satisfactory and has not generated spurious features. In §2.1 we describe the sources included in our large biomolecule survey. In §2.2 we describe the spectral parameters of our observations.

Table 2 lists the molecular parameters of all the 41 observed transitions in our survey. Column (1) lists the species; Column (2) lists the transition. Column (3) lists the calculated rest frequency and its 2 σ standard deviation. Column (4) lists the product of the line strength and the square of the relevant dipole moment; and column (5) lists the upper state

⁵For a complete description of the observations taken between 1999 October and 2000 April toward W75N, see Watson et al. (2002).

⁶A technical description of this can be found at http://astron.berkeley.edu/~plambeck/technical.html

energy level of each transition. Each correlator configuration included the full K structure of CH₃CN, the two strongest ($10_{*,10}$ - $9_{*,9}$ A & E) CH₃COOH transitions at 111.5 GHz, a HCOOH transition at 108 GHz, and finally several strong transitions of HCOOCH₃ and CH₃CH₂CN. The spectral windows containing these transitions had a bandwidth of 50 MHz and were divided into 128 channels for a spectral resolution of 0.39 MHz per channel. All data were combined, imaged, and self-calibrated using the MIRIAD software package (Sault, Teuben & Wright 1995). The 1 σ rms noise level of the final maps was used to set upper limits to the column densities of the undetected molecular species searched for in our survey.

2.1. Sources

Table 3 lists the three sources surveyed for methyl cyanide (CH₃CN), ethyl cyanide (CH₃CH₂CN), methanol (CH₃OH), formic acid (HCOOH), acetic acid (CH₃COOH) and methyl formate (HCOOCH₃) in column (1). Columns (2) and (3) list the pointing center for each observation. Columns (4) and (5) list the distance to each source from the Sun and the distance to each source from the Galactic center (Galactocentric)⁷, respectively. Finally, columns (6) and (7) list the kinetic temperature and mass of each star-forming region. Each source was selected based on the conclusions in Remijan et al. (2003), that regions containing detactable amounts of the large biomolecule CH₃COOH are HMCs 1) between 200 and 2000 M_{\odot} , 2) within 7 kpc of the Galactic center and 3) are not known to show chemical differentiation between O and N chemistry. We discuss each source in detail below.

2.1.1. G19.61-0.23

G19.61-0.23 (α =18^h27^m38.^s1; δ =-11°56′39″ [J2000.0]) is a well studied high mass star forming region. Some of the earliest observations of G19.61-0.23 were performed by Matthews et al. (1977) as they identified radio continuum emission toward regions containing OH maser sources. They interpreted the large amount of extended structure as containing several HII regions. This region has been extensively studied for OH masers (Garay, Reid & Moran 1985; Forster & Caswell 1989), H₂O masers (Forster & Caswell 1989; Comoretto et al. 1990; Hofner & Churchwell 1996) and CH₃OH masers (Kalenskii et al. 1994; Larionov et al. 1999). Recently, one of the most extensive studies of the structure and molecular environments of the G19.61-0.23 regions was performed by Garay et al. (1998). Using high resolution (2″-3″) VLA

⁷For a description of how Galactocentric distance was determined, see Remijan et al. (2003).

observations, they mapped the ionized gas of this region using the $H66\alpha$ recombination line and the 1.6 GHz continuum as well as the molecular gas using the (2,2) inversion transition of NH₃. The ionized gas observations indicated at least 5 distinct HII regions implying the presence of an OB association. The NH₃ observations showed molecular emission from three distinct clumps, one of which lies very near the center of the cluster of HII regions (Garay et al. 1998). Since these observations, the G19.61-0.23 region has been observed in other molecular species such as CS (Shirley et al. 2003; Wu & Evans 2003; Larionov et al. 1999), HCN (Wu & Evans 2003) and CO (Hofner et al. 2000). However, the largest molecular species observed toward this region is methyl acetylene (CH₃C₂H) (Fontani et al. 2002) with the IRAM 30 m and FCRAO 14 m radio telescopes. Using an angular size of 30" measured by Hofner et al. (2000) from C¹⁷O emission, Fontani et al. (2002) measured a temperature of $T_{rot} = 56(7)$ K and a total CH₃C₂H source averaged column density of $N_{tot} = 1.1(3) \times 10^{15}$ cm⁻². They concluded from the rotation temperature diagram of the G19.61-0.23 region that the molecular gas was very close to LTE and thus, the measured rotational temperature is a good estimate of the kinetic temperature of the molecular gas. However, this region has never been surveyed at high resolution for large nitrogen bearing species such as CH₃CH₂CN or oxygen bearing molecules such as HCOOCH₃ and CH₃COOH. Therefore, our observations will not only further constrain the detectability of CH₃COOH toward high mass HMCs but also obtain information on the N and O chemistry present toward G19.61-0.23.

2.1.2. G45.47+0.05

G45.47+0.05 (α =19^h12^m25^s.6; δ =11°09′26″ [J2000.0]) is a high mass star forming region that has been extensively studied for OH masers (Matthews et al. 1977; Forster & Caswell 1989), H₂O masers (Forster & Caswell 1989; Churchwell, Walmsley & Cesaroni 1990), and CH₃OH masers (Zheng & Ling 1997; Larionov 1999). Churchwell, Walmsley & Wood (1992) searched this region for emission from CS, ¹³CO and CH₃CN. Using an excitation temperature of 100 K, they derived a total ¹³CO column density of N_{tot}=3.6×10¹⁷ cm⁻² and a CS column density of N_{tot}=1.3×10¹³ cm⁻². Initially, CH₃CN was not detected in this region. However, Olmi, Cesaroni & Walmsley (1993) conducted a survey of NH₃ and CH₃CN using the IRAM 30 m radio telescope. From their observations, they were able to construct rotational temperature diagrams and measured an approximate kinetic temperature of T_k=44(1) K and a total CH₃CN beam averaged column density of N_{tot}=2.5(5)×10¹³ cm⁻². For NH₃, they measured a rotational temperature of T_{rot}=39 K and a total NH₃ beam averaged column density of N_{tot}=4.4×10¹⁴ cm⁻². One of the most extensive molecular line surveys done toward this source was performed by Hatchell et al. (1998) using the James Clerk Maxwell Radio Telescope (JCMT). Their survey included lines of C¹⁷O, SO, SO₂, CH₃OH, ¹³CH₃OH,

CH₃CCH, CH₃CN, CH₃CH₂CN, CH₃OCH₃, HCOOCH₃, HC₃N, OCS and CCH. From their survey, they detected lines of CH₃CN, CH₃C₂H, CH₃OH, C¹⁷O, CO and CCH and possibly HCOOCH₃ and HC₃N. However, they concluded from their survey that the hot core region toward G45.47+0.05 must be very small (<1") or low density (<10⁸ cm⁻³). Most recently, Fontani et al. (2002) observed transitions of CH₃C₂H and measured a source size of 35(8)" and from that, a rotational temperature of 39(7) K and a total source averaged CH₃C₂H column density of N_{tot}=5.0(3)×10¹⁴ cm⁻². Our observations investigated this region at high resolution (<10 arcsec²) for large biologically interesting molecules including HCOOH and CH₃COOH, in an attempt to detect any molecular emission associated with a compact HMC.

2.1.3. W75 N

W75N ($\alpha = 20^h 38^m 36^s.6$; $\delta = 42^o 37' 32''$ [J2000.0]) is a massive star forming region containing several UC HII regions that has been a target of many surveys since its detection by Westerhout (1958). Primarily this region has been extensively studied for OH masers (e.g. Palmer, Goss & Whiteoak 2004; Palmer, Goss & Devine 2003; Gasiprong, Cohen & Hutawarakorn 2002), H₂O masers (e.g. Valdettaro et al. 2002; Lekht & Krasnov 2000; Torrelles et al. 1997), CH₃OH masers (Voronkov, Austin & Sobolev 2002; Lekht et al. 2002; Minier, Conway & Booth 2001), and searches for H₂CO masers (Mehringer, Goss & Palmer 1995). There have been very few surveys of this source for large molecules. The most recent was performed by Kalenskii, Slysh & Val'tts (2002). In that survey, they searched for thermal CH₃OH lines as well as lines of CH₃CN, HC₃N, HCOOCH₃, and SO₂ with the IRAM 30 m radio telescope. They successfully detected the CH₃OH transitions but no emission from any of the other molecular species. However, Pankonin et al. (2001) successfully detected emission from CH₃CN using the Heinrich Hertz Telescope. From their observations, they measured a rotational temperature $T_{rot}=135(33)$ K and a total CH₃CN beam averaged column density of $N_{tot}=2.3(3)\times10^{14}$ cm⁻² within the 35" beam. More recently, Watson et al. (2002) followed up the observations of Pankonin et al. (2001) toward W75N using the BIMA Array. They concluded that the CH₃CN emission coming from the hot core region was optically thick and corrections need to be made in order to more accurately determine T_{rot} and N_{tot} . Remijan et al. (2004) made an extensive study of CH₃CN toward the hot core regions W51 e1/e2. In that analysis, they corrected for optical depth and the results from the rotational temperature diagram analysis closely matched the statistical equilibrium analysis. Furthermore, they found CH₃CN is a good tracer of large hard-to-detect species such as CH₃COOH. In this analysis, we re-examined the CH₃CN data from Watson et al. (2002) and investigated how well CH₃CN traces larger molecular species.

3. Results

The three sources in Table 3 were surveyed for CH₃COOH, as well as the selected transitions of other large molecular species listed in Table 2 (see §2.2). Figure 1 summarizes the CH₃COOH survey toward each region. The spectral line labels located at the upper left of each panel correspond to the rest frequency for the v_{LSR} of each source (40 km s⁻¹ for G19.61-0.23, 62 km s⁻¹ for G45.47+0.05, and 10 km s⁻¹ for W75N). The 1 σ noise level is shown at the left of each spectrum. No new sources of CH₃COOH were detected. We did make the first detections of HCOOH, HCOOCH₃ and CH₃CH₂CN toward G19.61-0.23 and HCOOH and HCOOCH₃ toward W75N. And finally, we detected CH₃CN emission from all the sources and were able to measure the column density of the detected molecular species and based on the 1 σ detection limit, place upper limits on the species that were not detected. In §3.1, we discuss the spectral lines detected toward G19.61-0.23 . In §3.2 and §3.3, we discuss the spectral lines detected toward G45.47+0.05 and W75N, respectively .

3.1. G19.61-0.23

Figure 2a shows the first detection of formic acid (HCOOH) toward G19.61-0.23. The spectral line label located in the upper left of panel a) corresponds to the rest frequency for a v_{LSR} of 40 km s⁻¹. The 1 σ noise level is shown at the left of the spectral window. Figure 2b shows the map of the 5_{1,4}-4_{1,4} transition of formic acid (contours) on top of the 108.1 GHz continuum emission (gray scale). Notice the emission peak of HCOOH is offset from the continuum emission. The offset in emission peaks is significant given the size of the synthesized beam (shown at the bottom left of Fig. 2b). Figure 3 shows the detection and distribution of several large molecular species toward the G19.61-0.23 region including CH₃OH, CH₃CH₂CN (EtCN) and HCOOCH₃ (MeF). Figures 3a and 3b show the detection of CH₃OH at 107.014 and 107.160 GHz, respectively. The spectral line labels are similar to Figure 2a. Figure 3f shows the distribution of the 107.014 GHz CH₃OH emission (contours) on the 3 mm continuum emission (gray scale) toward G19.61-0.23. Figures 3c-3e show the detections of several transitions of CH₃CH₂CN. Figure 3g shows the distribution of the 107.502 GHz CH₃CH₂CN emission. Finally, figures 3d and 3i show the detections of three transitions of HCOOCH₃. Figure 3h shows the distribution of the 107.537 GHz HCOOCH₃ emission. Figure 4 shows the detection and distribution of methyl cyanide (CH₃CN) toward the G19.61-0.23 region. In Figure 4f we show the predicted locations of the J=6-5 $K=0,1,2 \& 3 \text{ CH}_3^{13}\text{CN}$ emission lines. While the K=3 line is clearly seen near v_{LSR} of 148 km s⁻¹, the lower K value transitions are not as clearly seen. However, we note that as much as 50% of the emission from the K = 5 CH₃CN line may be due to the K = 0 &

1 CH₃¹³CN lines. Higher signal-to-noise observations are needed to confirm the amount of contamination from CH₃¹³CN. The peak positions of CH₃CN, CH₃OH, HCOOH, CH₃CH₂CN and HCOOCH₃ agree, strongly indicating that all the molecular emission is coming from a common source. Table 4 lists the molecular species that were detected toward G19.61-0.23 at or near the systemic velocity (40 km s⁻¹). Least square Gaussian fits were made for each spectral line in order to obtain the radial velocities and integrated line intensities for the detected transitions. In Table 4, column (1) lists either the calculated rest frequencies of previously identified species or the observed frequencies of unidentified transitions. Columns (2) and (3) list the molecular species and associated transitions, respectively. Finally, column (4) lists the fitted integrated line intensity of each detected transition.

3.2. G45.47+0.05

Figure 5 shows the results of our survey toward G45.47+0.05. Of the 7 molecular species searched for, we only detected emission from one transition of CH₃OH and 4 transitions of CH₃CN. Figure 5a shows the detection of the 3(0)-4(1) A⁺ transition of CH₃OH toward G45.47+0.05. The spectral line label located in the upper left of panel a) corresponds to the rest frequency for a v_{LSR} of 62 km s⁻¹. The 1 σ noise level is shown at the left of the spectral window. Figure 5b shows the map of the 3(0)-4(1) A⁺ transition of CH₃OH (contours) on top of the 107.0 GHz continuum emission (gray scale). In this case, significant differences between the continuum peak and the CH₃OH emission peak are not seen. Figures 5c-f show the average intensity maps and spectrum of the 4 detected transitions of the J=6-5CH₃CN emission lines. Once again, no offsets are seen between the CH₃CN emission peak and the 110.5 GHz continuum peak. However, from Figure 5c, there appears to be at least two additional emission peaks of CH₃CN, one to the southeast and one to the northwest. The peaks also appear in the K=1 window, however, we also see emission coming from the southwest. In the K=3 window, only the CH_3CN emission peak near the continuum is present. From the fitting routine to measure source sizes (§4.1), we found this region may contain several molecular cores. Higher signal-to-noise observations are necessary to identify the number of individual cores toward this region. Table 5 lists the molecular species that were detected toward G45.47+0.05 at or near the systemic velocity (62 km s⁻¹). The column labels are the same as in Table 4.

3.3. W75N

Figure 6 shows the first detections of HCOOCH₃ and HCOOH toward W75N. Figure 6a shows the detection of the 9_{2,8}-8_{2,7} A and E transitions of HCOOCH₃ at 107.5 GHz. The spectral line labels correspond to the rest frequency located at the upper left of the panel for a v_{LSR} of 10 km s⁻¹. The 1 σ noise level is shown at the left of the spectrum. Figure 6b shows the map of the 9_{2.8}-8_{2.7} A transition of HCOOCH₃ (contours) on top of the 107.5 GHz continuum emission (gray scale). Notice the emission peak of HCOOCH₃ is offset from the continuum emission. The offset in emission peaks is significant given the size of the synthesized beam (seen at the bottom left of Figure 6b). Figure 6c shows the detection of the 5_{1,5}-4_{1,4} transition of HCOOH at 108.1 GHz. This is the first time any large oxygen bearing molecule has been detected toward W75N. Furthermore, we do not detect any large nitrogen bearing molecules, such as CH₃CH₂CN, in this region. Figure 7 shows the average intensity maps and spectrum of the J=6-5 CH₃CN emission lines. From these maps, we see the CH₃CN emission is extended beyond the 4."4×2."7 synthesized beam (shown at the bottom left of each panel) but peaks near the continuum. The peaks of the emission contours of CH_3CN and $HCOOCH_3$ are offset by at least $\sim 2.$ "7, indicating that the molecular emission is coming from at least two distinct sources. Table 6 lists the molecular species that were detected toward W75N at or near the systemic velocity (10 km s⁻¹). The column labels are the same as in Table 4.

4. Discussion

4.1. Rotational Temperature Diagrams of CH₃CN

One of the most difficult factors to determine is the column densities of large biomolecules in the interstellar medium (ISM). By measuring the column densities of biomolecules in many different regions with various physical conditions, we can get a better understanding of their formation mechanism, distribution and relative abundance compared with other large molecular species.

There are three primary procedures used in calculating the total beam averaged column density of a particular molecular species. They are the rotational temperature diagram method (RTDM), the statistical equilibrium method (SEM) and to assume local thermodynamic equilibrium (LTE). The methods are outlined in detail in Remijan et al. (2004). With the exception of CH₃CN, not enough transitions with a wide spread of upper state energy levels have been observed to utilize the RTDM. Furthermore, because many of our species are asymmetric top molecules, detailed collision rates are not known so we cannot utilize the

SEM. In the case of CH₃CN, we can use the RTDM to determine a rotational temperature and the column density of CH₃CN. If the CH₃CN population can be described by a single rotational temperature and if all the lines are optically thin, the measured intensities are proportional to the column densities in the upper levels of the transitions being observed. From Remijan et al. (2004), we have

$$\frac{N_u}{g_u} = 2.04 \frac{\int \Delta I dv}{\Omega_s \nu^3 S_{ij} \mu^2} \times 10^{20} \text{cm}^{-2}, \tag{1}$$

and

$$\frac{N_u}{g_u} = \frac{N_{tot}}{Q(T_{rot})} e^{-(E_u/kT_{rot})},\tag{2}$$

where N_u is the upper-level column density and g_u is the statistical weight of the upper level, Ω_s is the solid angle subtended by the source (arcsec²), $\int \Delta I dv$ is the integrated line intensity (Jy beam⁻¹ km s⁻¹), ν is the rest frequency (GHz), and $S_{ij}\nu^2$ is the product of the total torsional-rotational line strength and the square of the electric dipole moment (Debye²) and in equation (2), N_{tot} is the total source averaged column density, $Q(T_{rot})$ is the rotational partition function, E_u is the upper state energy level, and T_{rot} is the rotational temperature. One can thereby determine the "rotational" temperature from the slope of the fit and an estimate of the total molecular column density from the y-intercept. In order to measure Ω_s and construct accurate rotational temperature diagrams, we need a CH₃CN source size for each region. Source sizes are determined for the detected transitions assuming a gaussian distribution and are summarized in Table 7. In column (1), we list the source. Columns (2) and (3) list the molecular species and associated transition, respectively. Finally, the last column lists the source size deconvolved from the synthesized beam determined from that transition along with its 1 σ error. From this analysis, we find an average source size of 43.3(8.4) arcsec² for CH₃CN, 2.6(2) arcsec² for CH₃CH₂CN, and 2.8(4) arcsec² for HCOOCH₃ toward G19.61-0.23. Toward G45.47+0.05, we find an average source size of 242.2(80.6) arcsec² for CH₃CN and 65.4(6.0) arcsec² for CH₃OH. Toward W75N, we find an average source size of 10.3(2.6) arcsec² for CH₃CN and 13.5(2.0) arcsec² for HCOOCH₃. Using these values, we constructed rotational temperature diagrams using CH₃CN for each source. The results are shown in Figure 8. In each diagram, optical depth corrections were made using the procedure outlined in Remijan et al. (2004).

Figure 8a shows the rotational temperature diagram (RTD) of CH₃CN toward G45.45+0.05. We were only able to detect 4 K-components in the J=6-5 ladder. From the RTD, we find a rotational temperature of 81(63) K and a total source averaged column density of

 $N_{tot}=1.9(2.3)\times10^{13}$ cm⁻². It is apparent from the spectrum shown in figure 5 that the CH_3CN emission is optically thin ($\tau < 0.003$). Given the lower rotational temperature and the highly extended (>240 arcsec²) nature of the emission it appears that CH₃CN may contain several individual cores that trace a cooler component of molecular gas. Figure 8b shows the RTD for all 6 K-components in the J=6-5 ladder toward W75N. From the diagram, we find a rotational temperature of 242(84) K and a total source averaged column density of $N_{tot}=7.5(3.9)\times10^{15}$ cm⁻². Furthermore, the emission is very compact (~10 arcsec²) however, the optical depth of these lines is no greater than 0.06. Toward W75N, we are clearly detecting a hot molecular core region unlike what is seen toward G45.47+0.05. Finally, figure 8c shows the RTD toward G19.61-0.23. In this case, we see the CH₃CN emission is optically thick and we obtain a non-physical rotational temperature. Therefore, the standard procedure of iteratively making corrections to the RTD cannot be applied in this case. However, if we assume $T_{rot}=230$ K (Kurtz et al. 2000), we can make optical depth corrections by changing the column density until the slope of the RTD gives a temperature of 230 K. Using this method, we find a total source averaged column density of $N_{tot}=1.6\times10^{18}$ ${\rm cm^{-2}}$ and optical depth values as large as 8.0. However, this value should only be considered an approximation and in order to truly measure the temperature, the column density of CH₃CN, and H₂ density toward this region, a full statistical equilibrium analysis needs to be applied.

4.2. Column Densities of Large Molecular Species in G19.61-0.23

In order to determine the column densities of the remaining large molecules, we use the rotational temperatures determined from CH₃CN and assume LTE excitation to determine the partition function for each species. The rotational partition functions for all the observed molecular species are given in Table 2. Figure 9 and Table 8 summarizes the total source averaged column densities, fractional abundances and upper limits of all the large molecular species toward each high mass star forming region. In Table 8, source for each column density measurement is given in the column (1). Columns (2), (3), (4) and (5) list the total source averaged column densities for HCOOH, HCOOCH₃, CH₃CH₂CN and CH₃COOH, respectively. We adopt the temperature reported by Kurtz et al. (2000) of 230 K for G19.61-0.23. Assuming this temperature we find an average total source averaged column density of N_{tot} =9.3×10¹⁶ cm⁻² for CH₃CH₂CN and N_{tot} =3.4×10¹⁷ cm⁻² for HCOOCH₃. For HCOOH, we find a total source averaged column density of N_{tot} =6.0×10¹⁶ cm⁻². Finally, using the 1 σ noise level of our combined data set, we were able to set an upper limit to the total column density of N_{tot} =3.0×10¹⁶ cm⁻² for CH₃COOH. This gives a relative CH₃COOH/HCOOCH₃ abundance ratio of less than 8.8×10⁻² for G19.61-

0.23. This is higher than the abundance ratio found toward Sgr B2(N-LMH), W51e2 and G34.3+0.2 (3-7×10⁻²) but this is still an upper limit (Remijan et al. 2003). The relative HCOOH/HCOOCH₃ abundance ratio is 0.18 which is comparable to the abundance ratios found by Liu et al. (2001) toward Sgr B2(N-LMH), Orion and W51 (\sim 0.10). For the two transitions of CH₃OH, we can use the same procedure to calculate an average CH₃OH total source averaged column density of $N_{tot}=2.0\times10^{18}$ cm⁻². However, both are maser transitions of CH₃OH so the column density determined from equation (1) is not valid in this case and should be considered only as an upper limit. From our observations, it appears that the hot core regions towards the G19.61-0.23 region are very similar to the hot cores in G34.3+0.2. In fact the measured column densities of HCOOCH₃ and CH₃CH₂CN are higher toward G19.61-0.23 than G34.3+0.2 (Remijan et al. 2003). We can conclude that the non-detection of CH₃COOH toward this region is signal-to-noise limited and that a deeper survey should detect emission of CH₃COOH. In fact, in Figure (1a), there is a suggestion that the 111.507 GHz line of CH₃COOH is starting to be resolved out of the noise.

4.3. Column Densities of Large Molecular Species in W75N and G45.47+0.05

In total we observed two transitions of HCOOCH₃, and one transition of HCOOH toward W75N and one transition of CH₃OH toward G45.47+0.05. For the one transition of CH₃OH toward G45.47+0.05 using the rotational temperature of 81 K from the RTDM of CH₃CN, we find a CH₃OH total source averaged column density of $N_{tot}=1.8\times10^{15}$ cm⁻². However, this is a maser transition of CH₃OH so the column density determined should once again only be considered an upper limit. For HCOOCH₃ toward W75N using the rotational temperature of 242 K from the RTDM of CH₃CN, we find an average total source averaged column density of N_{tot} =2.1×10¹⁶ cm⁻² and for HCOOH, we find a total source averaged column density of $N_{tot}=5.4\times10^{15}$ cm⁻². The relative HCOOH/HCOOCH₃ abundance ratio is 0.26 which is more than twice as large compared to the abundance ratios found by Liu et al. (2001). Figure 9 summarizes the column densities and upper limit column densities of CH₃COOH, CH₃CH₂CN, HCOOH, CH₃CN and HCOOCH₃ toward G19.61-0.23, G45.47+0.05 and W75N. Toward G45.47+0.05, we did not detect emission from any other large molecular species except for CH₃OH and CH₃CN. The emission from both species is highly extended: >60 arcsec² for CH₃OH and >240 arcsec² for CH₃CN and the temperature determined from the RTDM is only 81 K with a large error (63 K). Thus, it appears that G45.47+0.05, does not contain a compact hot molecular core that contains the physical conditions suitable to form large molecular species. In fact, the G45.47+0.05 region may contain a chemistry similar to cold dark clouds. Further observations that compare the chemistry of these regions are necessary to confirm this hypothesis.

W75N does contain hot cores but these observations suggest for the first time that there is a differentiation between the O and N cores toward this region. The emission from HCOOCH₃ and HCOOH is located ~2."7 to the east of the CH₃CN peak and has a cometary morphology similar to what was seen by Liu et al. (2003) toward Orion. The CH₃CN emission, on the other hand, falls very close to the continuum emission which is offset from the HCOOCH₃ and HCOOH emission. Indeed, the W75N region is more similar to the hot core regions surrounding W3 and Orion that show a differentiation between the O and N cores, than to the W51e2 or Sgr B2(N-LMH) HMCs. Further high resolution observations are necessary to investigate the O and N cores toward this region.

4.4. Implications for Future Acetic Acid (CH₃COOH) Surveys

Recalling the conclusions reached from Remijan et al. (2003), that CH₃COOH sources are HMCs located within 7 kpc of the Galactic center, had a mass range of 200 and 2000 M_{\odot} and did not show a differentiation between O and N bearing cores, we discuss the validity of these criteria in light of the new non-detections. The distances and mass requirements are met due to source selection (except for W75N, which at a galactocentric distance of 8 kpc is outside the 7 kpc limit. However, given the mass and size of this source and the errors involved in distance determination, we believed W75N would be a good source to survey for CH₃COOH). G45.47+0.05 we determined does not contain a hot compact core. The implications from the CH₃CN rotation diagram is that the molecular emission is extended and may contain several distinct sources. Because we believe CH₃COOH is formed in warm HMCs, we would not expect to detect it toward this region. We find that W75N has different O and N cores similar to $W3(OH)/W3(H_2O)$ or the Orion Compact Ridge/Hot Core. With this discovery, we would not expect to detect CH₃COOH in this region. We can also rule out the O and N differentiation is due to our ability to resolve out the individual O and N cores in these sources. While it is true that each source where O and N chemical differentiation has been seen is at a heliocentric distance of <3 kpc, we note that G19.61-0.23 does not show such a distinction at a heliocentric distance of 3.6 kpc. This leaves G19.61-0.23 as the lone remaining suitable CH₃COOH source. G19.61-0.23 seemed to fit all the selection criteria, and in fact, there is an indication of an CH₃COOH transition at 111.507 GHz, however, the corresponding A transition at 111.549 GHz is not seen (Figure 1). It appears that the chemistry present toward G19.61-0.23 is similar to G34.3+0.2. If you compare Figure 3 and Figure 3 from Remijan et al. (2003), it is seen that all the same transitions of each molecular species are detected; however, the spectral lines toward G19.61-0.23 are all weaker by a factor of $\sim 2-3$. If we apply the same reduction to the acetic acid lines, we find the acetic acid transitions should have a peak line intensity between 0.07 and 0.1 Jy beam⁻¹. The noise level in the acetic acid windows toward G19.61-0.23 was ~ 0.05 Jy beam⁻¹, so we would only expect to detect these lines at the 1-2 σ level. Thus, to be confident in a detection, we would need to decrease the noise level to ~ 0.02 Jy beam⁻¹ for a 3-5 σ detection. We expect further observations with better signal-to-noise should detect these lines. Therefore in summary, our new results remain consistent with the conclusions of Remijan et al. (2003). Apparently, the formation of CH₃COOH favors HMCs with well mixed N and O, despite the fact that CH₃COOH does not contain a N atom. If proved to be true this is an important constraint on CH₃COOH formation and possibly other structurally similar biomolecules.

CH₃CN appears to be a good proxy in the search for hard-to-detect large molecular species, particularly in light of the potential importance of N chemistry. This is because CH₃CN is easily detectable and has a column density similar to the hard-to-detect large molecules (Remijan et al. 2003). Figure 9 illustrates this toward each of our SFRs. While CH₃CN appears to have a column density closer to the large molecular species, it is still not clear how CH₃CN is related to the chemistry that forms these large molecules, especially in those regions that have O and N rich cores. The only analysis done trying to explain the formation of O and N rich cores was by Rodgers & Charnley (2001). They determined that the O and N rich cores are formed due to the abundance (or lack of) ammonia (NH₃) on grain mantles and the timescale of the core formation; O rich cores form first, then they become N rich. However, our observations indicate that large molecules are found in regions where the formation of O and N rich cores is occurring simultaneously in the interstellar environment. Because CH₃CN lines are easy to detect, by first observing a source for CH₃CN, one can obtain information about the column densities of the larger molecules and from that, an approximate line intensity. This will undoubtedly influence future surveys and at the same time get an understanding of the physical conditions present towards future surveyed sources.

5. Summary

We have surveyed three high mass Galactic star forming regions for interstellar methanol (CH₃OH), formic acid (HCOOH), acetic acid (CH₃COOH) methyl formate (HCOOCH₃), methyl cyanide (CH₃CN) and ethyl cyanide (CH₃CH₂CN) with the BIMA Array. From our observations, we have detected two new sources of interstellar HCOOH toward the hot core regions G19.61-0.23 and W75N. Assuming a temperature of 230 K for the G19.61-0.23 region, we find a HCOOH total source averaged column density of $N_{tot}=6.0\times10^{16}$ cm⁻². Using a temperature of 242 K based on the rotation diagram of CH₃CN toward W75N, we find a HCOOH total source averaged column density of $N_{tot}=5.4\times10^{15}$ cm⁻². We have also made the first detections of CH₃CH₂CN and HCOOCH₃ toward G19.61-0.23. For CH₃CH₂CN, we find an

average total source averaged column density of $N_{tot}=9.3\times10^{16}~cm^{-2}$ and for HCOOCH₃, we find an average total source averaged column density of $N_{tot}=3.4\times10^{17}~cm^{-2}$. The relative HCOOH/HCOOCH₃ abundance ratio is 0.18 which is comparable to the abundance ratios found by Liu et al. (2001) toward Sgr B2(N-LMH), Orion and W51 (\sim 0.10). We have made the first detection of $\mathrm{HCOOCH_3}$ toward W75N. For $\mathrm{HCOOCH_3}$, we find an average total source averaged column density of N_{tot}=2.1×10¹⁶ cm⁻². The relative HCOOH/HCOOCH₃ abundance ratio is 0.26 which is more than twice as large compared to the abundance ratios found by Liu et al. (2001). Furthermore, the hot core regions around W75N show for the first time a chemical differentiation between the O and N cores similar to what is seen toward the Orion Hot Core and Compact Ridge and W3(OH) and W3(H₂O). It is also apparent from our observations that the high mass star forming region G45.47+0.05 does not contain any compact hot molecular core and its chemistry may be more similar to cold dark clouds. Finally, it appears that CH₃CN is a good tracer for harder to detect biomolecules. CH₃CN was detected toward all three sources and its total source averaged column density was similar to the larger molecular species. Because CH₃CN lines appear easy to detect, by first observing a source for CH₃CN, one can obtain information about the column densities of the larger biomolecules and from that, an approximate line intensity. This will undoubtably influence future surveys and at the same time get an understanding of the physical conditions present towards future surveyed sources.

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Fig. 1.— (a-b) Acetic acid (CH₃COOH) survey lines toward G19.61-0.23. The noise level is \sim 60 mJy beam⁻¹ (indicated by the vertical bar at the left of each panel.). The spectral line labels correspond to the rest frequency located at the top left of each panel for a v_{LSR} of 40 km s⁻¹ (dashed line). c-d) same as a-b) except toward G45.47+0.05; rms \sim 40 mJy beam⁻¹; v_{LSR} =40 km s⁻¹. e-f) same as a-b) except toward W75N; rms \sim 50 mJy beam⁻¹; v_{LSR} =10 km s⁻¹.

Fig. 2.— (a) HCOOH spectrum toward G19.61-0.23. The noise level is \sim 40 mJy beam⁻¹ (indicated by the vertical bar at the left of the spectrum.). The spectral line label corresponds to the rest frequency located at the top left of the panel for a v_{LSR} of 40 km s⁻¹ (dashed line). (b) Emission contours from the $5_{1,5}$ - $4_{1,4}$ transition of HCOOH at 108.127 GHz overlaid on the G19.61-0.23 continuum emission at 108.1 GHz (gray scale). The numbers on the gray scale wedge are in units of Jy beam⁻¹. Contours indicate the location of the HCOOH emission near its peak at 40 km s⁻¹. The contour levels are -0.10, 0.10, 0.15, 0.20, 0.25, 0.30, and 0.35 Jy beam⁻¹. The synthesized beam of 5."5×2."9 is indicated at the bottom left corner.

Fig. 3.— (a-e) Spectrum taken toward G19.61-0.23 showing several detections of CH₃OH, CH₃CH₂CN (EtCN) and HCOOCH₃ (MeF). The rms noise level is indicated by the vertical bar at the left of the spectrum. The spectral line labels correspond to the rest frequency located at the top left of each panel for a v_{LSR} of 40 km s⁻¹ (dashed line). (f) Emission contours from the 3(0)-4(1) A⁺ transition of CH₃OH at 107.014 GHz at the emission peak of 40 km s⁻¹ overlaid on the G19.61-0.23 continuum emission at 107.0 GHz (gray scale). The numbers on the gray scale wedge are in units of Jy beam⁻¹. The CH₃OH contour levels are $-0.10, 0.10, 0.15, 0.20, 0.25, 0.30, \text{ and } 0.35 \text{ Jy beam}^{-1}$. The synthesized beam of $16.''9 \times 5.''5$ is indicated at the bottom left corner. (g) Emission contours from the 125,*-115,* transition of CH₃CH₂CN at 107.502 GHz at the emission peak of 40 km s⁻¹ overlaid on the G19.61-0.23 continuum emission at 107.5 GHz (gray scale). The CH₃CH₂CN contour levels are -0.1, 0.4, $0.6, 0.8, 1.0, 1.2, 1.4 \text{ and } 1.6 \text{ Jy beam}^{-1}$. The synthesized beam of $4.7^{\circ} \times 1.9^{\circ}$ is indicated at the bottom left corner. (h) Emission contours from the 9_{4,5}-8_{4,4} A transition of HCOOCH₃ at 111.453 GHz at the emission peak of 40 km s⁻¹ overlaid on the G19.61-0.23 continuum emission at 111.5 GHz (gray scale). The HCOOCH₃ contour levels are -0.10, 0.15, 0.20, 0.25, 0.30, 0.35, and 0.40 Jy beam⁻¹. The synthesized beam of $4.7^{\circ}\times1.7^{\circ}$ is indicated at the bottom left corner. (i) Spectrum showing the 111.453 GHz transition of HCOOCH₃. The labels are the same as in (a).

Fig. 4.— (a-d) Averaged intensity maps of the J = 6 - 5 CH₃CN emission lines. The gray scale in (a) is for the 110.4 GHz continuum of the G19.61-0.23 region. The numbers on the gray scale wedge are in units of Jy beam⁻¹. In each panel, the given K-ladder transition taken at the peak velocity of the line is mapped in contours. The contour levels are -0.10, 0.15,

0.20, 0.25, 0.30, 0.35 and 0.40 Jy beam⁻¹. The synthesized beam of 16."4×5."3 is indicated at the bottom left of each panel. (e) CH₃CN spectrum taken toward the G19.61-0.23 region. The rest frequency located in the top left corner corresponds to the K=0 component for a v_{LSR} of 40 km s⁻¹. The noise level is ~0.6 Jy beam⁻¹ (indicated by the vertical bar at the left of the spectrum.). (f) Same as (e) except the rest frequency corresponds to the K=4 component. Also shown are the predicted locations of the J=6-5 K=0,1,2&3 CH₃¹³CN emission lines. The K=3 line is clearly seen near v_{LSR} of 148 km s⁻¹.

Fig. 5.— (a) CH₃OH spectrum toward G45.47+0.05. The noise level is \sim 84 mJy beam⁻¹ (indicated by the vertical bar at the left of the spectrum.). The spectral line label corresponds to the rest frequency located at the top left of the panel for a v_{LSR} of 62 km s⁻¹ (dashed line). (b) Emission contours from the 3(0)-4(1) A⁺ transition of CH₃OH at 107.014 GHz overlaid on the G45.47+0.05 continuum emission at 107.0 GHz (gray scale). The numbers on the gray scale wedge are in units of Jy beam⁻¹. Contours indicate the location of the CH₃OH emission near its peak at 65 km s⁻¹. The contour levels are -0.15, 0.225, 0.30, 0.375, 0.45, and 0.525 Jy beam⁻¹. The synthesized beam of 9."6×6."1 is indicated at the bottom left corner. (c-f) Averaged intensity maps and spectrum of the J = 6 - 5 CH₃CN emission lines. The gray scale in (c) is for the 110.4 GHz continuum of the G45.47+0.05 region. The numbers on the gray scale wedge are in units of Jy beam⁻¹. In each panel, the given K-ladder transition taken at the peak velocity of the line is mapped in contours. The contour levels are -0.3, 0.3, 0.375, 0.45, 0.525, 0.60 and 0.75 Jy beam⁻¹. The synthesized beam of 9."3×5."8 is indicated at the bottom left of each panel.

Fig. 6.— (a) HCOOCH₃ spectrum toward W75N. The noise level is ~ 80 mJy beam⁻¹ (indicated by the vertical bar at the left of the spectrum.). The spectral line label corresponds to the rest frequency located at the top left of the panel for a v_{LSR} of 10 km s⁻¹ (dashed line). (b) Emission contours from the $9_{2,8}$ - $8_{2,7}$ E transition of HCOOCH₃ at 107.543 GHz overlaid on the W75N continuum emission at 107.5 GHz (gray scale). The numbers on the gray scale wedge are in units of Jy beam⁻¹. Contours indicate the location of the HCOOCH₃ emission near its peak at 10 km s⁻¹. The contour levels are -0.1, 0.1, 0.15, 0.20, 0.25 and 0.30 Jy beam⁻¹. The synthesized beam of 4."4×3."7 is indicated at the bottom left corner. (c) HCOOH spectrum toward W75N. The spectral line labels are similar as in (a).

Fig. 7.— Same as in figure 4, but toward W75N. The contour levels are 0.4, 0.5, 0.6, 0.7, 0.8 and 1.0 Jy beam⁻¹ except in d) where the contour levels are 0.2 and 0.3 Jy beam⁻¹. The synthesized beam of $2."7 \times 2."2$ is indicated at the bottom left corner of each panel.

Fig. 8.— (a) Rotational temperature diagram for the observed CII₃CN transitions toward G45.47+0.05. Ordinate: $\ln(N_u/g_u)$. Abscissa: $E_u(K)$. The least-squares fit across K-levels

is shown as a solid line. The resulting rotational temperature is 81 K. (b) Rotational temperature diagram for the observed CH₃CN transitions toward W75N. The resulting rotational temperature is 242 K. (c) Rotational temperature diagram for the observed CH₃CN transitions toward G19.61-0.23. In this case, the diagram gives a non-physical rotational temperature.

Fig. 9.— a) Beam-averaged column densities determined from G19.61-0.23 interferometric measurements. The unhatched bars indicate upper limits to the column density measurements; the hatched bars indicate measured column densities. b) Beam-averaged column densities determined from G45.47+0.05 interferometric measurements. c) Beam-averaged column densities determined from W75N interferometric measurements.

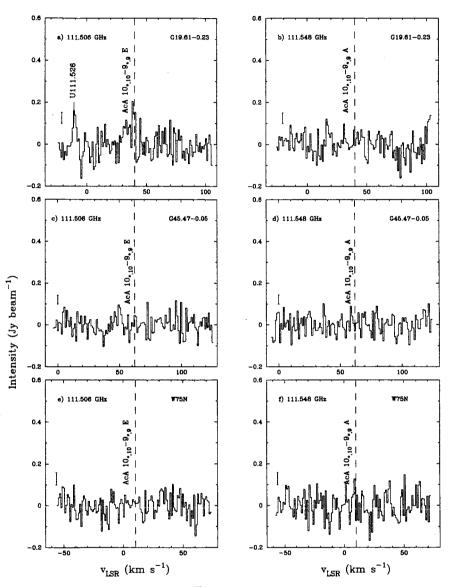


Figure 1.

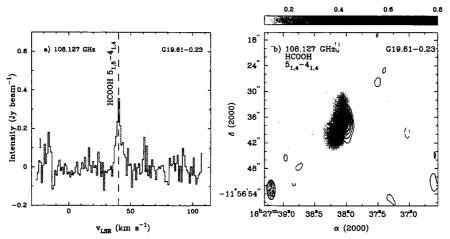


Figure 2.

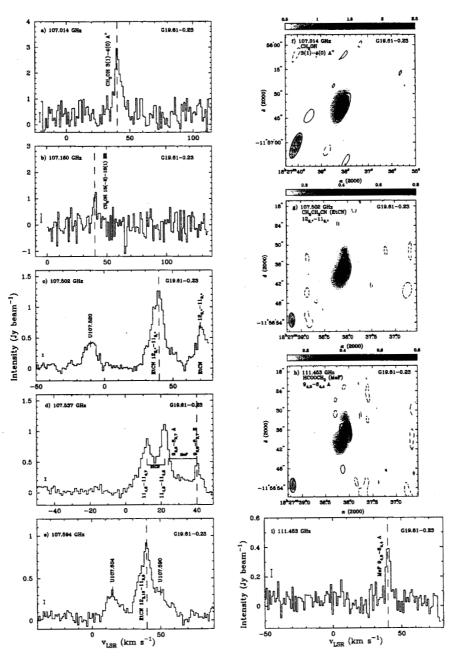
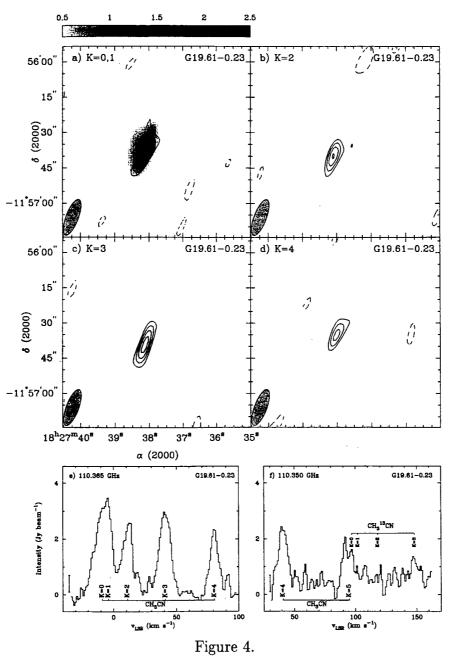


Figure 3.



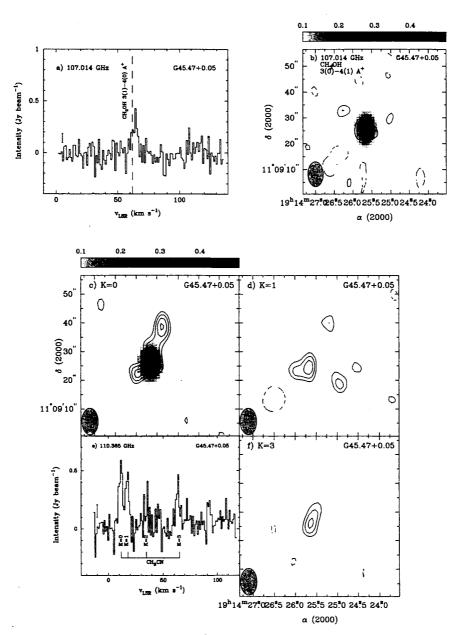


Figure 5.

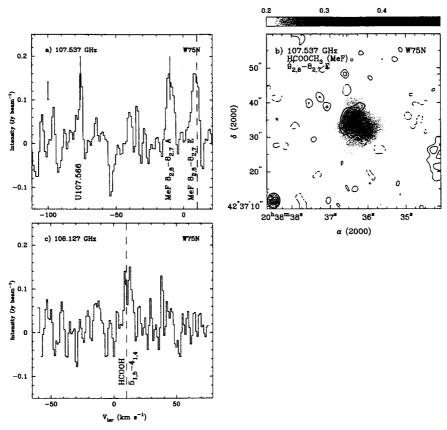


Figure 6.

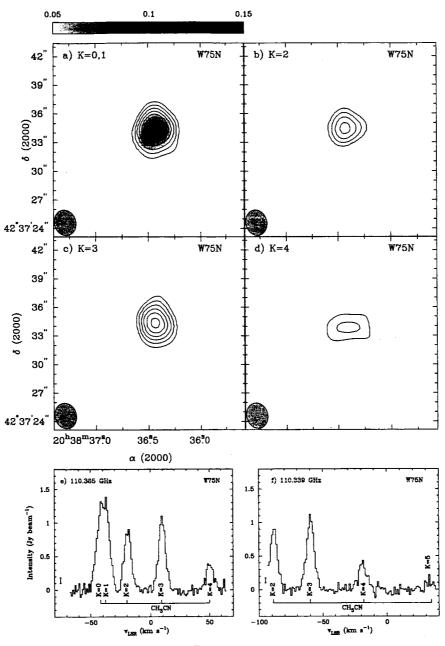


Figure 7.

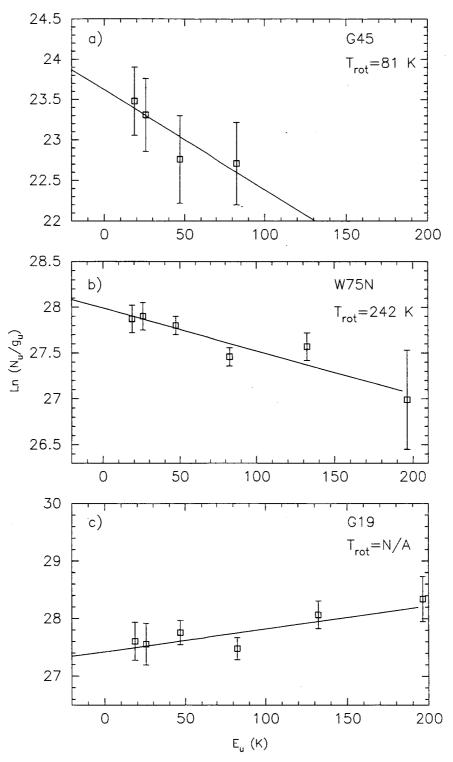


Figure 8.

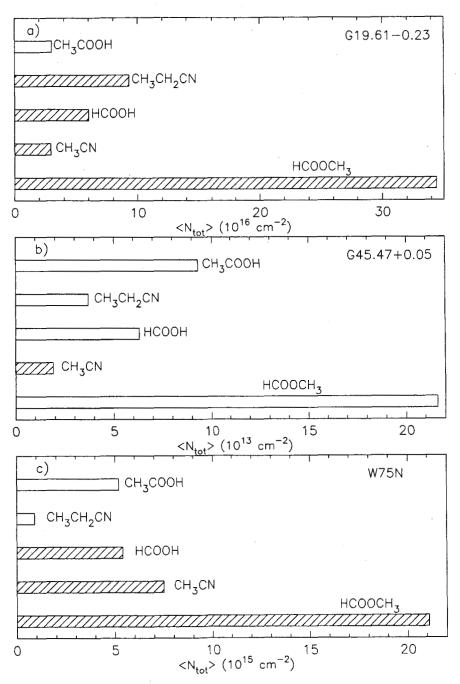


Figure 9.

Table 1. Combined Map Parameters

Source	Gain calibrator	beam ("×")	Frequency (GHz)	chan rms ^a (Jy beam ⁻¹)
G19.62-0.23	1733-130	4.5×1.9	111.5	0.05
		$16.9{\times}5.4$	110.3	0.60
G45.47+0.05	1925+211	3.3×1.9	111.5	0.04
		$9.3 { imes} 5.8$	110.3	0.10
W75Nb	2025+337	2.7×2.2	110.3	0.07
	2013+370	4.4×3.7	111.5	0.06

^aThe stated channel rms is given for the window containing the transitions of AcA at 111.5 GHz or the CH₃CN transitions at 110.3 GHz.

 $^{^{\}rm b} Data$ toward W75N at 110.3 GHz used with permission from C. Watson (2003Oct16).

Table 2. Molecular Line Parameters

Species	Transition	Frequency ^a (MHz)	$< S_{i,j} \mu^2 > $ (Debye ²)	$\mathbf{E}_{m{u}}$ (K)
$\mathrm{CH_3CN^b}$	6(5)-5(5)	110,329.608(120)	28.1	196.30
	6(4)-5(4)	110,349.760(200)	51.0	132.34
	6(3)-5(3)	110,364.470(100)	68.9	82.57
	6(2)-5(2)	110,374.986(120)	81.6	47.02
	6(1)-5(1)	110,381.376(120)	89.3	25.68
	6(0)-5(0)	110,383.494(120)	91.9	18.56
CH ₃ ¹³ CN ^b	6(3)-5(3)	110,309.847(10)	68.9	83.0
	6(2)-5(2)	110,320.438(10)	81.6	47.2
	6(1)-5(1)	110,326.795(10)	89.3	25.7
	6(0)-5(0)	110,328.914(10)	91.9	18.5
CH ₃ OH ^c	3(1)-4(0) A+	107,013.850(100)	3.01	28.4
	15(-2)-15(1) E2	107,159.790(100)	2.61	305.0
CH ₃ OD ^d	3(1)-3(0) E1	110,475.760(1826)	3.8	22.2
	4(1)-4(0) E1	110,950.750(2382)	5.1	30.8
CH ₃ COOH ^e	10 _{*,10} -9 _{*,9} E	111,507.270(40)	54.8	30.5
	10 _{*,10} -9 _{*,9} A	111,548.533(40)	54.8	30.5
CH ₃ CH ₂ CN ^f	172,16-171,17	107,481.468(46)	91.5	70.0
	$12_{7,6}$ - $11_{7,5}$	107,485.178(26)	117.3	88.1
	$12_{7,5}-11_{7,4}$	107,485.178(26)	117.3	88.1
	$12_{6,7}$ - $11_{6,6}$	107,486.961(26)	133.4	73.6
	$12_{6,6}$ - $11_{6,5}$	107,486.962(26)	133.4	73.6
	$12_{8,5}$ - $11_{8,4}$	107,491.573(28)	98.8	104.7
	$12_{8,4}$ - $11_{8,3}$	107,491.573(28)	98.8	104.7
	$12_{5,8}\text{-}11_{5,7}$	107,502.426(26)	147.0	61.4
	$12_{5,7}$ - $11_{5,6}$	$107,\!502.474(26)$	147.0	61.4
	$12_{4,9}$ - $11_{4,8}$	107,543.926(26)	158.1	51.4
	$12_{4,8}$ - $11_{4,7}$	107,547.601(26)	158.1	51.4
	123,10-113,9	107,594.049(28)	166.8	43.6
HCOOHg	51,5-41,4	108,126.710(160)	9.3	18.8
HCOOCH ₃ h	9 _{2,8} -8 _{2,7} E	107,537.189(24)	22.8	28.8

Table 2—Continued

Species	Transition	Frequency ^a (MHz)	$< S_{i,j} \mu^2 > $ (Debye ²)	Е _и (К)	
	9 _{2,8} -8 _{2,7} A	107,543.746(26)	22.8	28.8	
	9 _{8,2} -8 _{8,1} A	110,455.358(30)	69.1	5.1	
	9 _{8,1} -8 _{8,0} A	110,455.358(30)	69.1	5.1	
	$9_{8,2}$ - $8_{8,1}$ E	110,457.971(28)	69.1	5.1	
	$9_{7,2}$ - $8_{7,1}$ E	110,525.598(32)	59.2	9.6	
	9 _{7,3} -8 _{7,2} A	110,535.182(28)	59.2	9.6	
	9 _{7,2} -8 _{7,1} A	110,535.184(28)	59.2	9.6	
	$9_{7,3}$ - $8_{7,2}$ E	110,535.955(26)	59.2	9.6	
	$7_{2,6}$ - $6_{1,5}$ E	110,550.217(28)	19.0	1.4	
	7 _{2,6} -6 _{1,5} A	110,560.053(30)	19.0	1.4	
	9 _{4,5} -8 _{4,4} A	111,453.327(24)	19.4	37.3	

 $^{^{\}rm a}{\rm Errors}$ are 2 σ

Table 3. Hot Cores and Their Physical Parameters

Source Name	R.A. J(2000)	DEC. J(2000)	Distance (kpc)	Galactocentrica (kpc)	Τ _ε (K)	$\begin{array}{c} \text{Mass} \\ (\text{M}_{\odot}) \end{array}$
G19.61-0.23 ^b	18:27:38.1	-11:56:39	3.5	4.8	230	450
G45.47+0.05b	19:14:25.6	11:09:26	8.3	6.3	90	250
W75N	20:38:36.6	42:37:32	2.0^{c}	8.0	135 ^d	\sim 440 $^{\rm d}$

 $^{^{\}mathrm{a}}$ Galactocentric distance was determined using the procedure outlined in Remijan et al. (2003).

^bBoucher et al. (1980); $Q_{rot}=0.98T_{rot}^{3/2}$

^cXu & Lovas (1996); $Q_{rot}=1.23T_{rot}^{3/2}$

^dAnderson et al. (1988); Q_{rot} =0.71 $T_{rot}^{3/2}$

 $^{^{\}mathrm{e}}$ Illuyshin et al. (2001); Q $_{rot} \! = \! 14.10 T_{rot}^{3/2}$

fLovas (1982); $Q_{rot}=7.17T_{rot}^{3/2}$

 $^{^{\}rm g}$ Willemot et al. (1980); Q $_{rot} \! = \! 1.71 T_{rot}^{3/2}$

 $^{^{\}rm h}{\rm Oesterling}$ et al. (1999); ${\rm Q}_{rot}{=}12.45T_{rot}^{3/2}$

^bKurtz et al. (2000)

^cHunter et al. (1994)

^cPankonin et al. (2001)

cWatson et al. (2002)

Table 4. G19.61-0.23 Molecular Line Identifications

Frequency ^a (MHz)	Species	Transition	$\int \! \Delta I dv^b $ (Jy beam ⁻¹ km s ⁻¹)
107,013.850(100)	CH ₃ OH	3(1)-4(0) A+	24.1(3.9)
107,159.790(100)	CH ₃ OH	15(-2)-15(1) E2	4.5(1.7)
107,491.573(28)	$EtCN^c$	$12_{8,*}-11_{8,*}$	7.8(6)
107,502.426(26)	EtCN	$12_{5,*}-11_{5,*}$	14.7(8)
107,520.0	U107.520	, ,	5.3(5)
107,537.189(24)	MeF ^d	$9_{2,8}$ - $8_{2,7}$ E	3.1(3)
107,543.746(26)	MeF	9 _{2,8} -8 _{2,7} A	$3.1(1.0)^{e}$
107,543.926(26)	EtCN	124,9-114,8	5.5(1.0)e
107,547.601(26)	EtCN	124,8-114,7	6.6(4)
107,590.0	U107.590		3.2(5)e
107,594.049(28)	EtCN	$12_{3,10}$ - $11_{3,9}$	10.0(6)e
107,604.0	U107.604		3.9(4)
108,126.710(160)	HCOOH	$5_{1,5}$ - $4_{1,4}$	1.5(3)
110,309.847(10)	$\mathrm{CH_{3}^{13}CN}$	6(3)-5(3)	13.1(3.0)
110,329.608(120)	CH ₃ CN	6(5)-5(5)	16.3(2.1)
110,349.760(200)	CH ₃ CN	6(4)-5(4)	22.5(2.0)
110,364.470(100)	CH_3CN	6(3)-5(3)	33.7(2.3)
110,374.986(120)	CH_3CN	6(2)-5(2)	26.4(2.1)
110,381.376(120)	CH_3CN	6(1)-5(1)	24.0(3.0)
110,383.494(120)	CH_3CN	6(0)-5(0)	25.4(3.0)
111,453.327(24)	MeF	9 _{4,5} -8 _{4,4} A	1.8(3)
111,526.0	U111.526		0.5(2)

^aErrors are 2 σ

 $^{^{\}rm b}{\rm Errors}$ are 1 σ

 $^{^{}c}EtCN=Ethyl$ Cyanide (CH3CH2CN)

 $^{^{\}rm d}$ MeF=Methyl Formate (HCOOCH₃)

^eBlended line

Table 5. G45.47+0.05 Molecular Line Identifications

Frequency ^a (MHz)	Species	Transition	$\int \Delta I dv^b$ (Jy beam ⁻¹ km s ⁻¹)
107,013.850(100)	CH ₃ OH	3(1)-4(0) A ⁺	1.69(7)
110,364.470(100)	CH ₃ CN	6(3)-5(3)	1.6(3)
110,374.986(120)	CH_3CN	6(2)-5(2)	0.9(3)
110,381.376(120)	CH_3CN	6(1)-5(1)	1.9(3)
110,383.494(120)	CH ₃ CN	6(0)-5(0)	2.3(4)

 $^{^{\}rm a} \rm Errors$ are 2 σ

Table 6. W75N Molecular Line Identifications

Frequency ^a (MHz)	Species	Transition	$\int \! \Delta \mathrm{Idv^b} $ (Jy beam ⁻¹ km s ⁻¹)
107,537.189(24)	MeF ^c	9 _{2,8} -8 _{2,7} E	0.8(1) ^d
107,543.746(26)	MeF	9 _{2,8} -8 _{2,7} A	$0.7(1)^{d}$
107,566.0	U107.566	-,,	$0.5(1)^{d}$
108,126.710(160)	HCOOH	$5_{1,5}$ - $4_{1,4}$	$0.6(1)^{d}$
110,329.608(120)	CH ₃ CN	6(5)-5(5)	0.8(2)e
110,349.760(200)	CH ₃ CN	6(4)-5(4)	2.9(2)
110,364.470(100)	CH ₃ CN	6(3)-5(3)	7.6(3)
110,374.986(120)	CH ₃ CN	6(2)-5(2)	6.4(2)
110,381.376(120)	CH ₃ CN	6(1)-5(1)	7.6(4)e
110,383.494(120)	CH ₃ CN	6(0)-5(0)	$7.6(4)^{e}$

^aErrors are 2 σ

 $^{^{\}rm b} \rm Errors$ are 1 σ

 $^{^{\}mathrm{b}}\mathrm{Errors}$ are 1 σ

^cMeF=Methyl Formate (HCOOCH₃)

^dThe intensity and line width were approximated because the least square Gaussian fitting routine did not give a satisfactory fit.

^eBlended line

Table 7. Source Size Measurements

Source	Species	Transition	Size (arcsec ²)
G19.61-0.23	CH ₃ CN	6(0)-5(0)	60.9(6.0) ^a
	CH ₃ CN	6(1)-5(1)	$60.9(6.0)^{a}$
	CH ₃ CN	6(2)-5(2)	31.8(3.2)
	CH ₃ CN	6(3)-5(3)	15.5(3.0)
	CH ₃ CN	6(4)-5(4)	64.8(15.1)
	EtCN	$12_{8,*}-11_{8,*}$	2.9(2)
	EtCN	$12_{5,*}-11_{5,*}$	1.7(3)
	EtCN	$12_{4,8}$ - $11_{4,7}$	3.6(2)
	EtCN	$12_{3,10}$ - $11_{3,9}$	2.6(1)
	EtCN	124,9-114,8	$2.2(1)^{a}$
	MeF	9 _{2,8} -8 _{2,7} A	$2.2(1)^{a}$
	MeF	$9_{2,8}$ - $8_{2,7}$ E	2.4(5)
	MeF	9 _{4,5} -8 _{4,4} A	3.8(6)
	нсоон	$5_{1,5}$ - $4_{1,4}$	point source
	CH_3OH	$3(1)-4(0) A^+$	point source
G45.47+0.05	CH_3CN	6(0)-5(0)	243.1(35.5)
	CH_3CN	6(1)-5(1)	294.2(151.2)
	$\mathrm{CH_{3}CN}$	6(3)-5(3)	189.4(43.4)
	CH ₃ OH	$3(1)-4(0) A^+$	65.4(6.0)
W75N	CH_3CN	6(0)-5(0)	$11.3(3)^{a}$
	CH_3CN	6(1)-5(1)	$11.3(3)^{a}$
	CH_3CN	6(2)-5(2)	8.8(3)
	CH_3CN	6(3)-5(3)	8.8(3)
	CH_3CN	6(4)-5(4)	10.1(1.1)
	CH ₃ CN	6(4)-5(4)	12.7(5.0)
	MeF	9(2,8)-8(2,7) A	19.8(2)
	MeF	9(2,8)-8(2,7) E	7.2(2)

^aBlended line

Table 8. Column Densities and Fractional Abundances^a

Source	$N_{ m HCOOH} \ (10^{15} \ { m cm}^{-2})$	$N_{\rm HCOOCH_3} = (10^{15} {\rm cm}^{-2})$	${ m ^{N_{CH_{3}CH_{2}CN}}} (10^{15}~{ m cm^{-2}})$	$N_{\rm CH_3COOH} = (10^{15} {\rm cm}^{-2})$
G19.61-0.23	60.2	344.5	93.3	<29.6
G45.45+0.05	< 0.06	< 0.22	< 0.04	< 0.09
W75N	5.4	21.1	< 0.92	< 5.17
	X _{нсоон} (×10 ⁻⁸)	X _{HCOOCH3} (×10 ⁻⁸)	$X_{\text{CH}_3\text{CH}_2\text{CN}} \times 10^{-8}$	$X_{\text{CH}_3\text{COOH}} \times 10^{-8}$
G19.61-0.23 ^b	7.5	43.1	11.7	>3.7
G45.45+0.05b	>0.01	>0.04	>0.007	>0.02
W75N ^c	2.7	10.6	>0.5	>2.6

^aAll upper limits were calculated from the 1σ rms noise in the spectral window, and the source size. For the upper limit of HCOOCH₃, we used a frequency of 111.453 GHz, for CH₃CO₂CN, 107.492 GHz, for HCOOH, 108.126 GHz and for CH₃COOH, 111.549 GHz.

^bH₂ column density taken from Fontani et al. (2002)

^cH₂ column density taken from Zinchenko, Henkel & Mao (2000)